

Solving Schrodinger Equations Using Kashuri and Fundo Transform Decomposition Method

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Abstract

This paper proposes a hybrid method that combines the Kashuri and Fundo integral transform and the Adomian decomposition method to solve two types of Schrödinger equations. The exact solutions are successfully found, and the results are compared with those obtained by the existing methods. The obtained results show the accuracy and efficiency of the method.

1. Introduction

Differential equations (DEs) play a major role in describing a wide range of natural phenomena undergoing change. They relate an unknown function to its derivatives. Partial differential equations (PDEs), as well as, ordinary differential equations (ODEs) usually arise in many diverse applications, such as fluid flow, mechanical systems, physics, and relativity [1]. There are a variety of numerical and analytical techniques for solving the PDEs.

The analytical solution for a given PDE is always preferable because of the restriction of the numerical solutions which may not offer us much about the behavior of systems. However, in many instances, systems described by PDEs are frequently so intricate or so large where a purely analytical solution to the equations is not tractable, or many of the modeled problems are resulted in nonlinear PDEs in which the exact solutions are difficult to obtain by using analytical methods. In this case, the numerical methods are needed to be developed for solving DEs despite these methods have some errors.

Many numerical methods were developed over a period of time [2-7], in order to find more accurate approximation solutions to these nonlinear equations. On the other hand, predict the future behavior of a dynamic system play a major role in various fields of applied sciences such as plasma physics, hydrodynamic, nonlinear optics, and quantum mechanics. The exact solutions of these DEs are very important.

Schrödinger equations (SchEs) play a crucial role in numerous spaces of practical sciences [8] and the references therein. In classical mechanics, the Schrödinger linear

partial differential equation (SchLPDE) is a wave equation that plays an important role of Newton's second law where the future behavior of a dynamic system can be predicted. The wave function predicts the path of a given dynamical system or the distribution of results over time. The SchLPDE is given by:

$$u_t(x, t) = i u_{xx}, i = \sqrt{-1}, t \geq 0 \quad \dots(1)$$

with:

$$u(x, 0) = h(x) \quad \dots(2)$$

where $h(x)$ is a continuous and square integrable function. On the other hand, the non-linear Schrödinger equation (NLSchE) is given by:

$$i u_t + u_{xx} + \rho |u|^2 u = 0 \quad \dots(3)$$

with:

$$u(x, 0) = h(x) \quad \dots(4)$$

such that ρ is a constant term and u is complex.

Several different numerical schemes have been used for solving SchLPDEs and NLSchEs [9-14], and so on.

The main objective of this topic is to investigate the use of Kashuri-Fundo Decomposition Method (KF-DM) for solving SchLPDE and NLSchE. Using this method, we found the exact solutions of some applications successfully. The results are compared with the those obtained by other methods in literatures. The findings demonstrate the method's dependability, precision, and adaptability.

2. Preliminaries

In this part, we will look at the Kashuri and Fundo transforms and their definitions and characteristics. by [15,16]:

Definition 2.1. The Kashuri and Fundo transformation of the function $k(t) \in H$ where t is a positive real number, is defined by:

$$\mathbb{K}[k(t)](v) = A(v) = \frac{1}{v} \int_0^\infty e^{\frac{-t}{v^2}} f(t) dt, \quad -r_1 < v < r_2$$

where $\mathbb{K}[k(t)]$ denotes the Kashuri-Fundo transformation of the function $k(t)$, and $H = \{k(t) : \text{there exists positive real numbers } M, r_1, \text{ and } r_2 \text{ such that } |k(t)| < M e^{\frac{|t|}{r_1^2}}, \text{ if } t \in (-1)^j \times [0, \infty)\}$

We give some useful properties and basic theorems of Kashuri-Fundo transform.

Theorem 2.2 (Duality relation). Let $k(t) \in H$ with Laplace transform $F(s)$. The Kashuri-Fundo $A(v)$ of $k(t)$ is given by:

$$A(v) = \frac{1}{v} F\left(\frac{1}{v^2}\right) \quad \dots(5)$$

Theorem 2.3 (Linearity of Kashuri-Fundo Transform).

Let $k(t)$, $k_1(t)$ and $k_2(t)$ be functions whose a new integral transforms exists for $|v| < r$ and let c be a constant. Then for $|v| < r$,

1. $\mathbb{K}[(k_1 + k_2)(t)] = \mathbb{K}[k_1(t)] + \mathbb{K}[k_2(t)]$.
2. $\mathbb{K}[(ck)(t)] = c\mathbb{K}[k(t)]$.

Theorem 2.4 (Some fundamental properties of Kashuri-Fundo transform). Let $A(v)$ be a Kashuri-Fundo transform of $k(t)$. Then:

1. $\mathbb{K}[k'(t)] = \frac{A(v)}{v^2} - \frac{k(0)}{v}$.
2. $\mathbb{K}[k''(t)] = \frac{A(v)}{v^4} - \frac{k(0)}{v^3} - \frac{k'(0)}{v}$.
3. $\mathbb{K}[k^{(n)}(t)] = \frac{A(v)}{v^{2n}} - \sum_{j=0}^{n-1} \frac{k^{(j)}(0)}{v^{2(n-j)-1}}$.

Table 1. List of some special Kashuri-Fundo transformation of $k(t)$, [17].

$k(t)$	$\mathbb{K}[k(t)]$
1	v
$t^n, n \geq 0$	$n! v^{2n+1}$
$e^{-\mu t}$	$\frac{v}{1+\mu v^2}$
$\sin(\mu t)$	$\frac{a v^2}{1+\mu^2 v^4}$
$\cos(\mu t)$	$\frac{v^2}{1+\mu^2 v^4}$

3. Kashuri-Fundo Decomposition Method (KF-DM)

In this section, we provide the KF-DM to standard NLSchE (3)-(4). We start by insertion the Kashuri-Fundo transform on each side of Eq. (3):

$$\frac{1}{v^2} \mathbb{K}[u(x, t)] - \frac{1}{v} u(x, 0) - i\mathbb{K}[u_{xx}] - i\mathbb{K}[\rho|u|^2 u] = 0 \quad \dots(6)$$

and substituting Eq. (4) into Eq. (6), we get:

$$\mathbb{K}[u(x, t)] = v k(x) + i v^2 \mathbb{K}[u_{xx} + [\rho|u|^2 u]] \quad \dots(7)$$

Using the inverse of Kashuri-Fundo transformation to Eq. (7):

$$u(x, t) = k(x) + \mathbb{K}^{-1} [i v^2 \mathbb{K} [u_{xx} + \rho|u|^2 u]] \quad \dots(8)$$

where $u(x, t)$ is the unknown function and $k(x)$ is the source terms.

We will now suppose an infinite series solution for the function $u(x, t)$ of the form:

$$u(x, t) = \sum_{n=0}^{\infty} u_n(x, t) \quad \dots(9)$$

Note that $|v|^2 v = v^2 \bar{v}$ and it can be written as follows:

$$|u|^2 \bar{u} = \sum_{n=0}^{\infty} A_n \quad \dots(10)$$

where A_n are Adomian polynomials, which can be computed using the following formula:

$$A_n = \frac{1}{n!} \frac{d^n}{dx^n} [F(\sum_{r=0}^n \omega^r u_r)]_r \quad \dots(11)$$

where n is a nonnegative integer.

Some of A_n 's are:

$$\begin{aligned} A_0 &= F(u_0) \\ &= u_0^2 \bar{u}_0 \\ A_1 &= u_1 F'(u_0) \\ &= 2u_0 u_1 \bar{u}_0 + u_0^2 \bar{u}_1 \\ A_2 &= u_2 F'(u_0) + \frac{1}{2!} u_1^2 F''(u_0) \\ &= 2u_0 u_2 \bar{u}_0 + u_1^2 \bar{u}_0 + 2u_0 u_1 \bar{u}_1 + u_0^2 \bar{u}_2 \end{aligned}$$

Putting Eq. (9) and Eq. (10) into Eq. (8), yields:

$$\sum_{n=0}^{\infty} u_n(x, t) = k(x) + i\mathbb{K}^{-1} \left[v^2 \mathbb{K} \left[\sum_{n=0}^{\infty} u_{nxx} + \sum_{n=0}^{\infty} A_n \right] \right] \quad \dots(12)$$

The following recursive relation can be generated through comparing each sides of Eq. (12):

$$\begin{aligned} u_0(x, t) &= k(x) \\ u_1(x, t) &= i\mathbb{K}^{-1} [v^2 \mathbb{K} [u_{0xx} + A_0]] \\ u_2(x, t) &= i\mathbb{K}^{-1} [v^2 \mathbb{K} [u_{1xx} + A_1]] \\ u_3(x, t) &= i\mathbb{K}^{-1} [v^2 \mathbb{K} [u_{2xx} + A_2]] \end{aligned}$$

We can write the general recursive relation of the unknown function $u(x, t)$ as:

$$u_{n+1}(x, t) = i\mathbb{K}^{-1} [v^2 \mathbb{K} [u_{nxx} + A_n]] \quad \dots(13)$$

where n is a nonnegative positive integer number.

Solutions of $u(x, t)$ may be given exactly or approximately as:

$$u(x, t) = \sum_{n=0}^{\infty} u_n(x, t) \quad \dots(14)$$

4. Applications

In this section, the KF-DM will be applied to solve some SchLPDEs and NLSchEs.

Example 4.1. Consider the SchLPDE:

$$i u_t + u_{xx} = 0 \quad \dots(15)$$

with the initial condition given by:

$$u(x, 0) = \cosh(3x) \quad \dots(16)$$

Applying the KF-DM to Eq. (15) and using the initial condition given in Eq. (16), we get the recurrence relation given by:

$$\begin{aligned} u_0(x, t) &= \cosh(3x) \\ u_{n+1}(x, t) &= i\mathbb{K}^{-1} [v^2 \mathbb{K} [u_{nxx}]], \quad n \geq 0 \end{aligned} \quad \dots(17)$$

We express few components as:

$$u_0(x, t) = \cosh(3x) \quad \dots(18)$$

$$u_1(x, t) = i\mathbb{K}^{-1} [v^2 \mathbb{K} [u_{0xx}]] \quad \dots(19)$$

$$= i\mathbb{K}^{-1} [9 v^3 \cosh(3x)]$$

$$u_2(x, t) = iK^{-1}[v^2K[u_{1xx}]] \dots(20)$$

$$= iK^{-1}[81iv^5 \cosh(3x)]$$

$$= \frac{(9it)^2}{2!} \cosh(3x)$$

$$u_3(x, t) = iK^{-1}[v^2K[u_{2xx}]] \dots(21)$$

$$= iK^{-1}\left[\frac{277}{2!} i^2 v^7 \cosh(3x)\right]$$

$$= \frac{(9it)^3}{3!} \cosh(3x)$$

and so on. Thus, on summing the above iterations, we obtain:

$$u(x, t) = \cosh(3x) \left[9it + \frac{(9it)^2}{2!} + \frac{(9it)^3}{3!} + \dots\right] \dots(22)$$

which is leading to the exact solution:

$$u(x, t) = \cosh(3x) e^{9ix} \dots(23)$$

This agrees with the result found by (ADM) [18].

Example 4.2. Consider the following NLSchE of the form:

$$iu_t + u_{xx} + 6|u|^2u = 0 \dots(24)$$

with the initial condition given by:

$$u(x, t) = e^{3ix} \dots(25)$$

Applying the Kashuri-Fundo transform coupled with ADM:

$$u_0(x, t) = u(x, 0) \dots(26)$$

$$u_{n+1}(x, t) = iK^{-1}[v^2K(u_{nxx} + 6A_n)] \dots(27)$$

where:

$$A_0 = \overline{u_0} u_0^2 = e^{3ix}$$

$$A_1 = 2u_0 u_1 \overline{u_0} + u_0^2 \overline{u_1} = -3it e^{3ix}$$

$$A_3 = 2 u_0 u_2 \overline{u_0} + u_1^2 \overline{u_0} + 2u_0 u_1 \overline{u_1} + u_0^2 \overline{u_2}$$

$$= \frac{9}{2} i^2 t^2 e^{3ix}$$

$$\vdots$$

and so on.

$$u_0(x, t) = e^{3ix} \dots(28)$$

$$u_1(x, t) = iK^{-1}[v^2K(u_{0xx} + 6A_0)] \dots(29)$$

$$= iK^{-1}[v^2K(-3e^{3ix})]$$

$$= iK^{-1}[-3v^3 e^{3ix}]$$

$$= -3it e^{3ix}$$

$$u_2(x, t) = iK^{-1}[v^2K(u_{1xx} + 6A_1)] \dots(30)$$

$$= iK^{-1}[v^2K(9ite^{3ix})]$$

$$= iK^{-1}[9v^5 ie^{3ix}]$$

$$= \frac{9}{2!} i^2 t^2 e^{3ix}$$

$$= \frac{(3it)^2}{2!} e^{3ix}$$

$$u_3(x, t) = iK^{-1}[v^2K(u_{2xx} + 6A_2)] \dots(31)$$

$$= iK^{-1}\left[v^2K\left\{\frac{-27}{2} i^2 t^2 e^{3ix}\right\}\right]$$

$$= iK^{-1}[-27 i^2 v^7 e^{3ix}]$$

$$= \frac{-27}{3!} i^3 t^3 e^{3ix}$$

$$= -\frac{(3it)^3}{3!} e^{3ix}$$

$$\vdots$$

Then:

$$u(x, t) = \sum_{n=0}^{\infty} u_n(x, t) \dots(32)$$

$$= e^{3ix} \left(1 - 3it + \frac{(3it)^2}{2!} - \frac{(3it)^3}{3!} + \dots\right)$$

$$= e^{3i(x-t)} \dots(33)$$

Then the exact solution is:

$$u(x, t) = e^{3i(x-t)} \dots(34)$$

The same exact agrees with the result found by (ADM) [9] and (VIM) [10,11].

Example 4.3. Consider the following NLSchE of the form:

$$iu_t + u_{xx} + 2|u|^2u = 0 \dots(35)$$

with the initial condition given by:

$$u(x, 0) = e^{ix} \dots(36)$$

Applying the Kashuri-Fundo transform coupled with ADM:

$$u_0(x, t) = u(x, 0) \dots(37)$$

$$u_{n+1}(x, t) = iK^{-1}[v^2K(u_{nxx} + 2A_n)] \dots(28)$$

where:

$$A_0 = u_0^2 \overline{u_0} = e^{ix},$$

$$A_1 = 2u_0 u_1 \overline{u_0} + u_0^2 \overline{u_1} = i e^{ix},$$

$$A_2 = 2u_0 u_2 \overline{u_0} + u_1^2 \overline{u_0} + 2u_0 u_1 \overline{u_1} + u_0^2 \overline{u_2} = i^2 e^{ix},$$

$$\vdots$$

and so on.

$$u_0(x, t) = e^{ix} \dots(39)$$

$$u_1(x, t) = iK^{-1}[v^2K(u_{0xx} + 2A_0)] \dots(40)$$

$$= iK^{-1}[v^3 e^{ix}]$$

$$= i t e^{ix}$$

$$u_2(x, t) = iK^{-1}[v^2K(u_{1xx} + 2A_1)] \dots(41)$$

$$= iK^{-1}[i v^5 e^{ix}]$$

$$= \frac{1}{2!} t^2 i^2 e^{ix}$$

$$u_3(x, t) = iK^{-1}[v^2K(u_{2xx} + 2A_2)] \dots(42)$$

$$= iK^{-1}[v^7 e^{ix}]$$

$$= \frac{1}{3!} i^3 t^3 e^{ix}$$

Then:

$$u_0(x, t) = \sum_{n=0}^{\infty} u_n(x, t) \dots(43)$$

$$= e^{ix} \left(1 + it + \frac{i^2}{2!} t^2 + \frac{i^3}{3!} t^3 + \dots\right)$$

which is leading to the exact solution:

$$u(x, t) = e^{i(x+t)} \dots(44)$$

The same exact which agrees with the result found by (ADM) [17].

5. Conclusion

Our goal of this work is to provide a combined computational method consisting of an integral transform and a semi-analytical method for solving SchLPDEs and NLSchEs. It is an application of the Kashuri-Fundo transform with the Adomain method. The exact solutions of some SchLPDEs and NLSchEs were successfully found and the obtained results are compared to ADM and VIM. This approach shows its efficiency, accuracy, and simplicity. This approach can be used to solve various types of differential equations in physics, engineering, and applied mathematics.

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Conflicts of Interest

The authors declare that there is no conflict of interest.

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